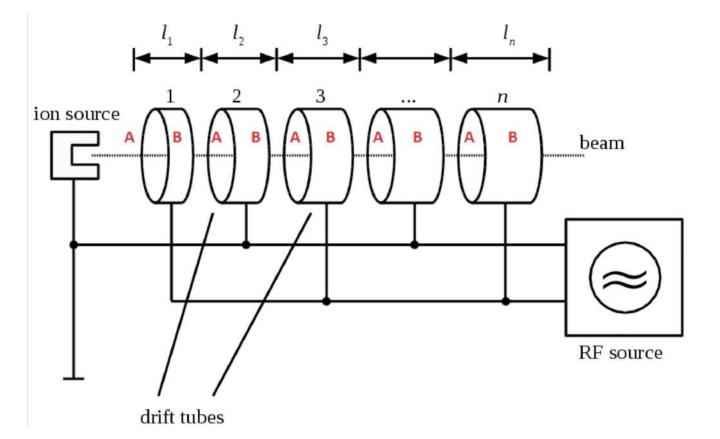
A Quick Look at Laser Wakefield Accelerators

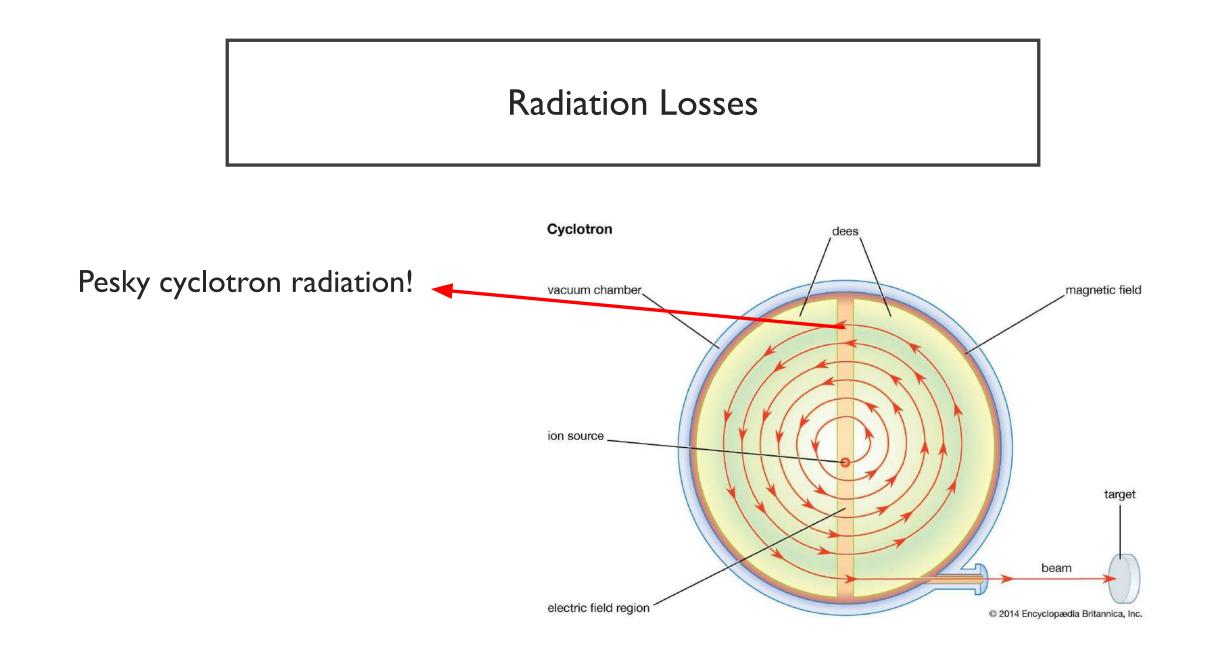
Bryan Dinh

Existing Accelerator Issues

Ionization Limits

- Require high voltages to achieve reasonable acceleration gradient
- Limited to E fields of 100 MV/m
- Needs to be long



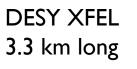


GeV Electron Sources and FELs

SLAC 3.2 km long

- Long wait time
- Expensive
- Inaccessible
- Large



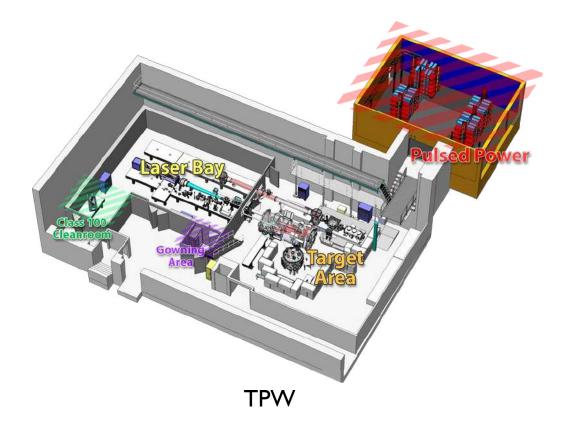


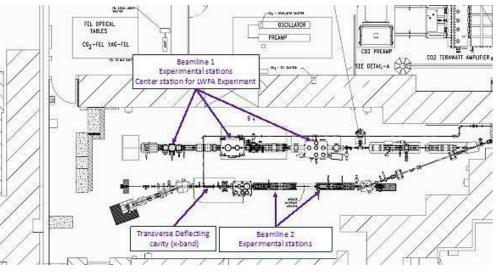


LWFA size

requirements

Room size only needs to on the order of magnitude of the size of a laser!





BNLATF

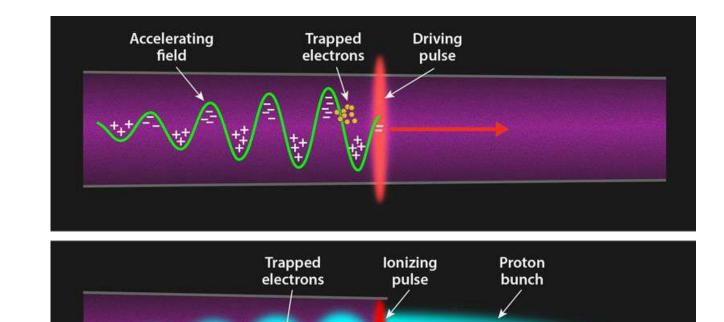


WHAT LWFA IS TACKLING

- High energy accelerator inaccessibility
- Skyrocketing cost of experiments
- Accelerator size
- High brightness-low charge applications

Laser Wakefield Accelerators (LWFA)

Plasma wakefield accelerators



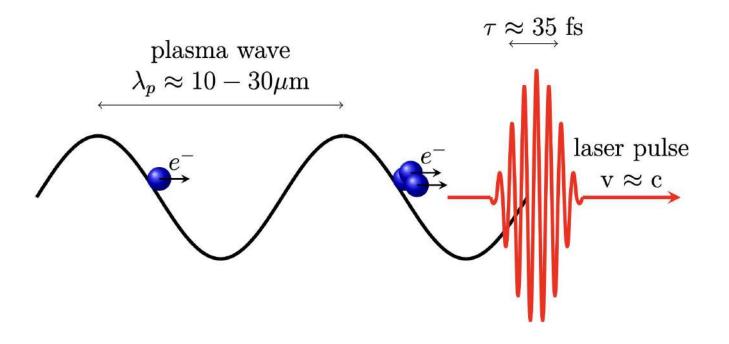
Laser Driver

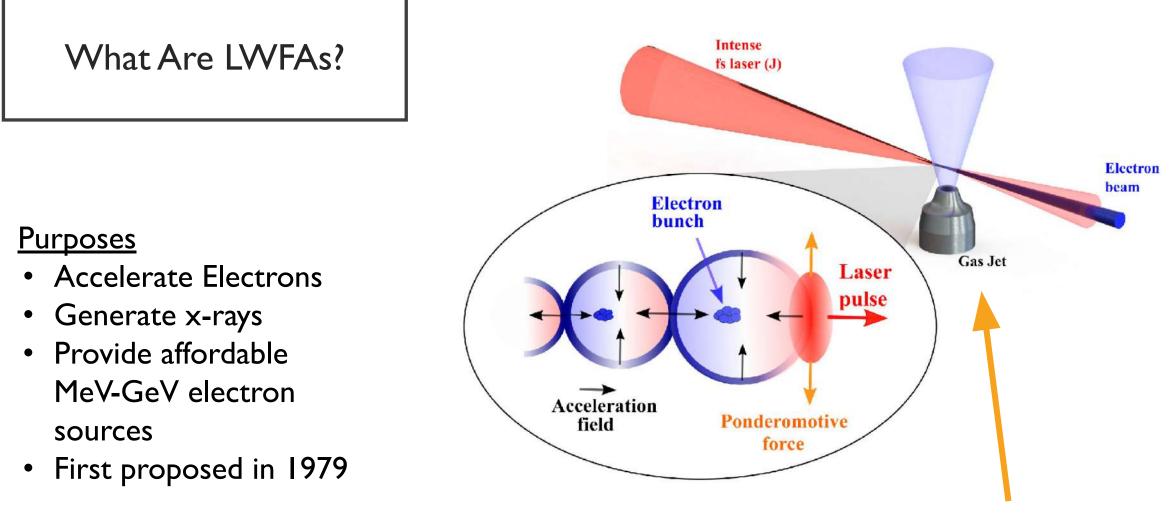
Proton Driver

What Are LWFAs?

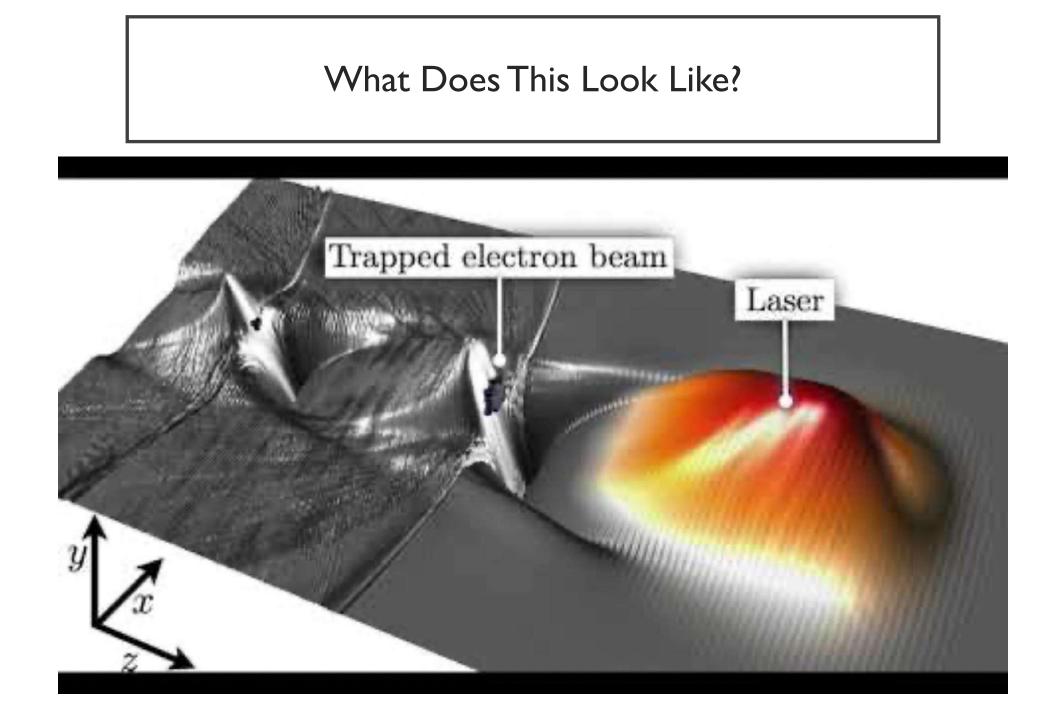
<u>Purposes</u>

- Accelerate Electrons
- Generate x-rays
- Provide affordable MeV-GeV electron sources
- First proposed in 1979





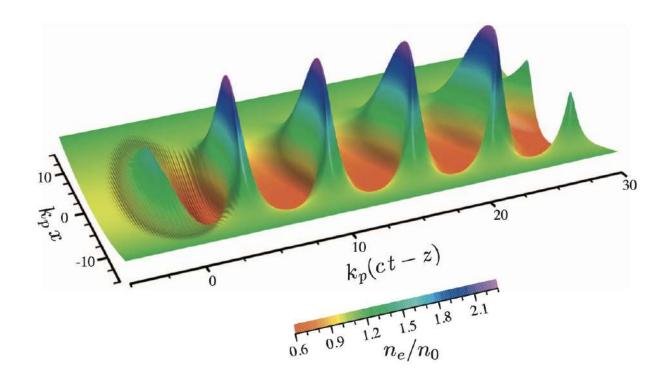
Gas jet size is around ~ 1-2 cm E > TeV/m



LWFA Basics

LASER-PLASMA INTERACTION

- The laser's field perturbs the plasma
- Creates plasma wake



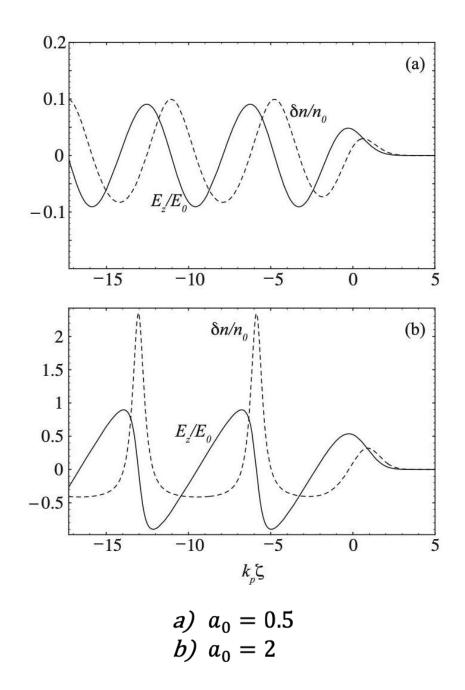
$$\omega_p = \sqrt{\frac{n_e e^2}{m_e \varepsilon_0}}$$



As laser strength increases, the plasma waves experience wave-breaking

 $E_{max} = m_e \omega_p v_{phase} / e$

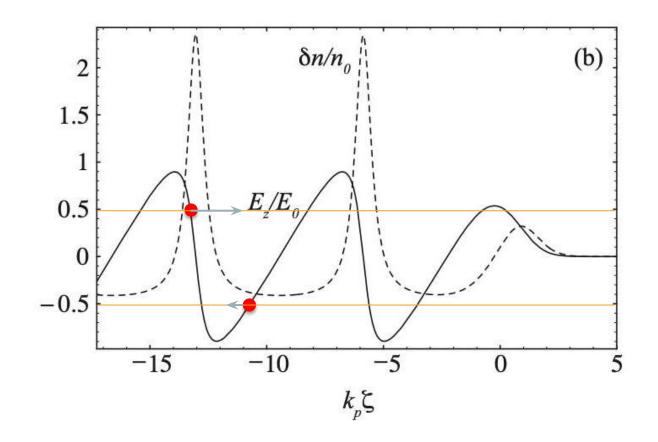
This occurs because the work done by the electric becomes comparable to the relativistic mass



Ponderomotive Force

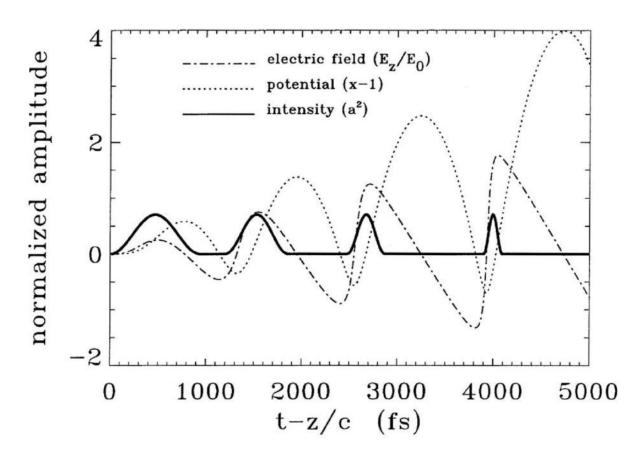
$$F_p = -m_e c^2 \nabla(\frac{a^2}{2})$$

- Requires the steepest gradients possible.
- We need a high-powered laser!
- Requires nonlinear interactions!



Beat Wave and Multi-Pulse LWFA

- Constituted the earliest LWFA experiments
- Needed when only low powered lasers were available
- Multiple pulses fire consecutively to resonantly excite plasma wakes



Self-modulated LWFA

- When laser is strong enough to drive nonlinear wakes in a single pulse
- Not strong enough to completely void space of electrons
- Laser is modulated by nonlinear interactions with the plasma

Forward Raman Instability

- Scattering occurs when laser strongly interacts with the plasma
- This is sign of resonant energy transfer
- Laser-plasma is said to be coupled

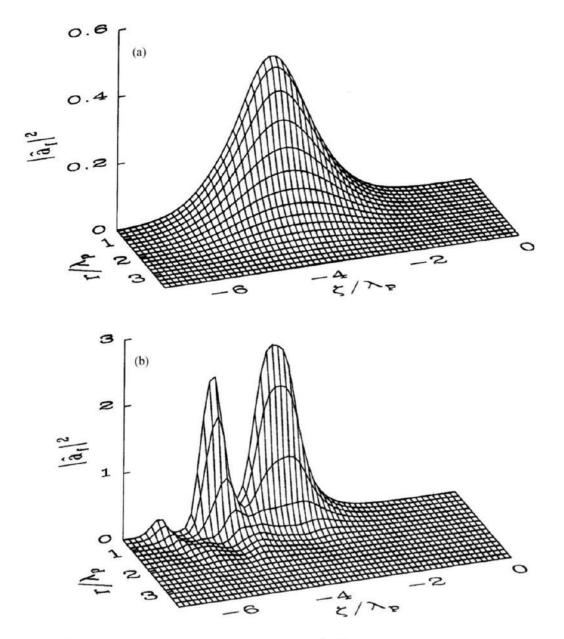


FIG. 15. Normalized laser intensity $|a|^2$ for the self-modulated LWFA case at (a) $c\tau=2Z_R$ and (b) $c\tau=3.2Z_R$. Laser pulse is moving to the right. From Krall *et al.*, 1993

Self-modulated LWFA

- As laser modulates, some electrons spend more time in the accelerating stage than others
- Terrible energy distribution

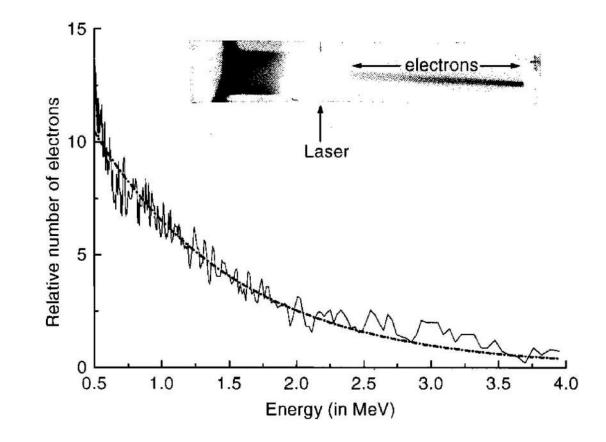


FIG. 9. Energy spectrum of accelerated electrons using direct exposure x-ray film. The dashed curve is an exponential fit to the data.

CHIRPED PULSE AMPLIFICATION (CPA)

- First demonstrated in 1985
- Allows the generation of ultrapowerful and ultrashort beams
- An intensity of ~ $2E18 W/cm^2$ required for nonlinear effects

MAINE et al.: GENERATION OF PEAK POWER PULSES BY CPA

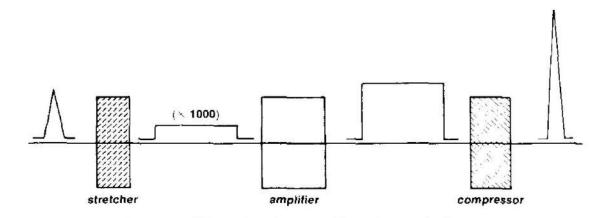
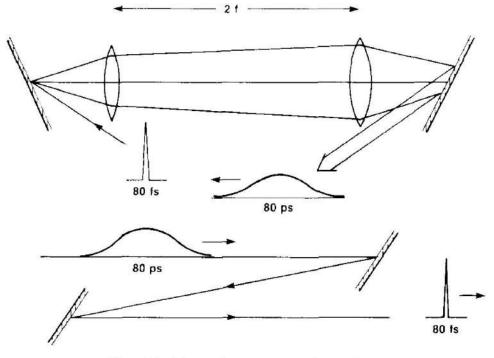


Fig. 1. Chirped pulse amplification technique.

Stretcher and Compressor Layout



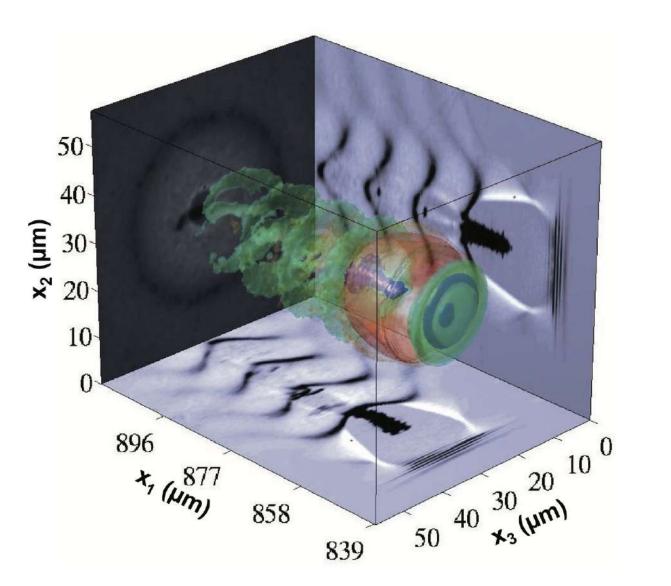
Can now readily obtain TW class, fs pulsed lasers

 UT3's 30 TW laser can produce >80 MeV electrons

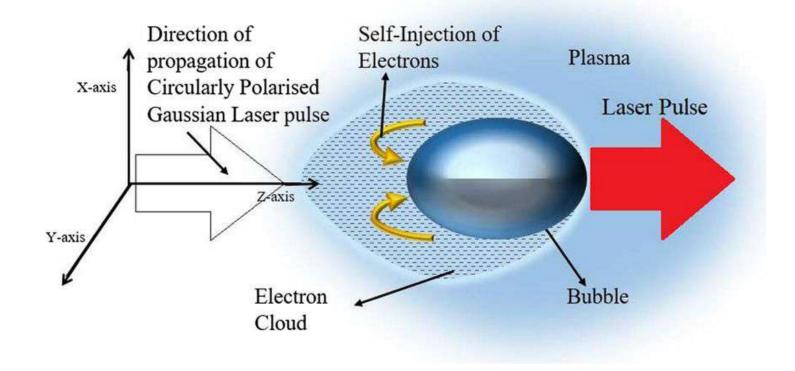
Fig. 11. Expansion-compression setup.

Bubble Regime LWFA

- Laser is powerful enough to create a bubble completely void of electrons in the initial wake period
- Forms quasi-monoenergetic distributions
- First demonstrated in 2004

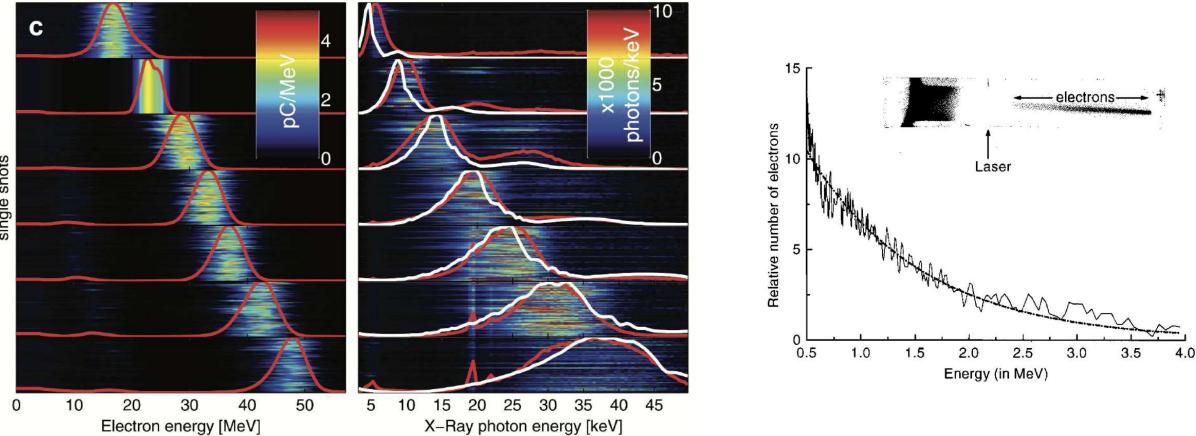






Bubble Regime

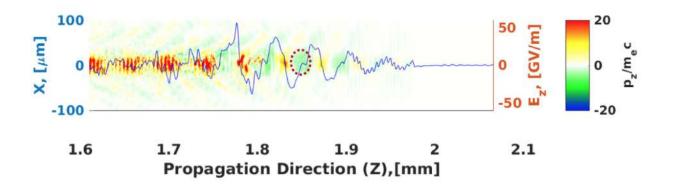
Self-Modulated Regime



single shots

Electron Injection

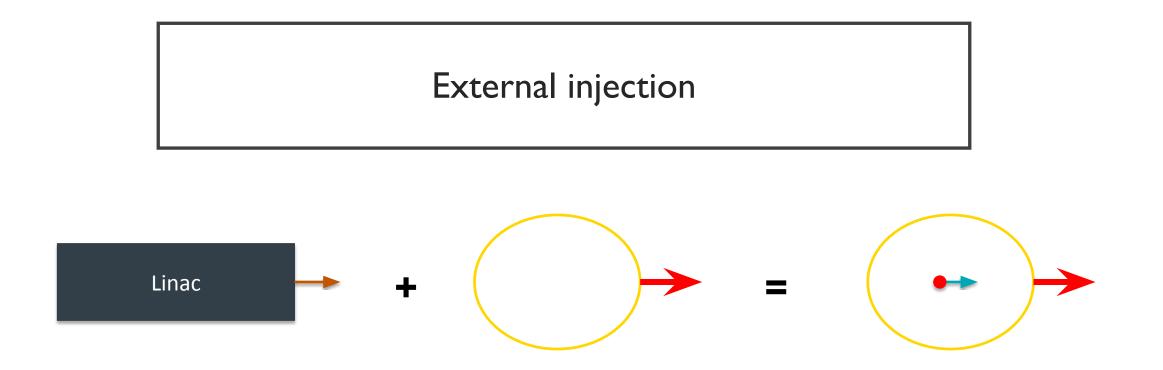




- Electrons that undergo DLA gain longitudinal momentum before interacting with nonlinear fields
- Supported by electrons with frequencies matching the laser frequency

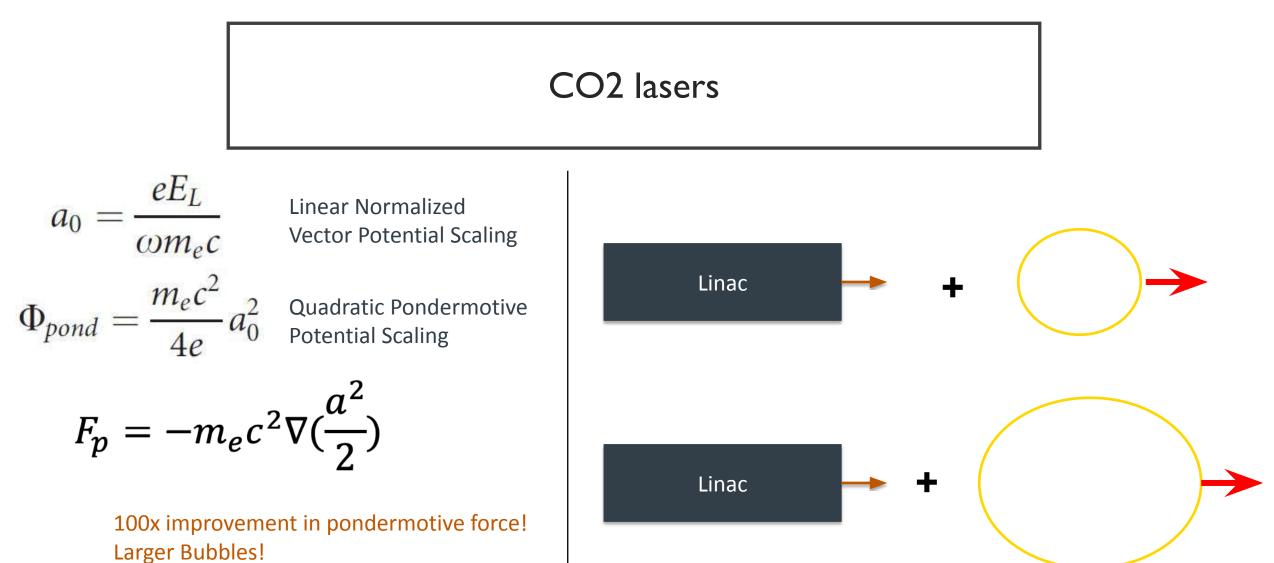


- Inject a small number of heavier atoms into gas mix to alter when ionization occurs
- Nitrogen is commonly used (1%-10%)



Developing injection scheme

- Traditional LWFA experiments use drive lasers ~
 I micron
- New CO2 lasers can now produce TW, fs pulses at 10 microns



Pulse Guiding

Kerr Effect: Nonlinear Self-Foc

 $\frac{\delta n}{n_0}$ FOCUSING DIFFRACTING a an rst $2\lambda_p$ 0 λ_p

- Longer laser-plasma interaction times offer greater attainable energies
- Diffraction must be dealt with

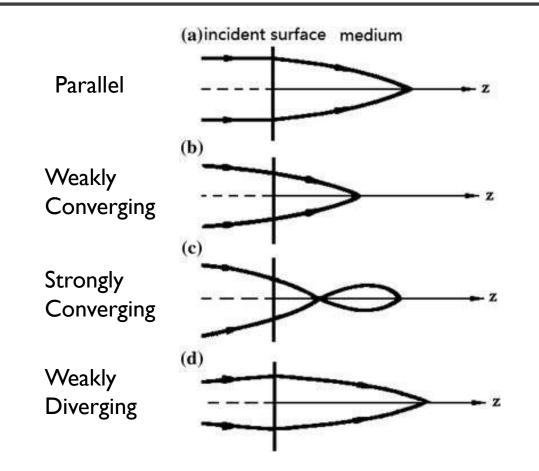
 $P^{(3)}(t) = \frac{1}{4} \varepsilon_0 \chi^{(3)} E^3 \operatorname{co:}_{\text{ated plasma wave on an initially uniform low-intensity}}^{\mathrm{FIG. 34. Schematic of focusing effects of an externally}_{\mathrm{pulse.}}$

 $n = n_0 + n_2 I$

E(t) =

The laser's intensity focuses itself!

Kerr Effect: Nonlinear Self-Focusing



Plasma Channeling

As light goes from a low to high index of refraction, it bends towards the normal of the interface

Methods:

- I. Pre-pulse
 - I. Transverse Axial Focusing Probe
 - 2. Longitudinal pre-pulse
- 2. Special gas jets

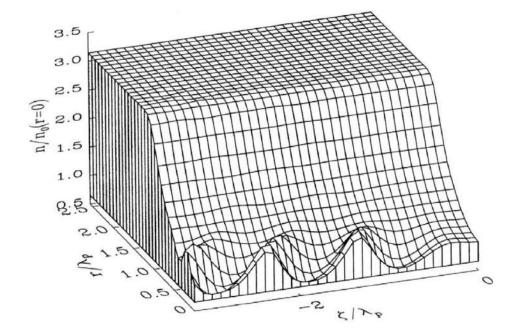
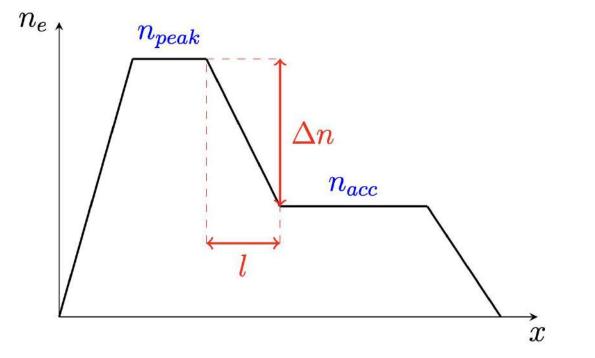


FIG. 29. Plasma electron density n/n_0 at $c\tau=20Z_R$ for a channel-guided LWFA. Initial density profile is parabolic with a depth $\Delta n = \Delta n_c = 1/\pi r_e r_0^2$. From Esarey, Sprangle, *et al.*, 1993.

Gas Jet Density Profile

- Higher density allows for early collection of charge
- Lower density allows for less laser diffraction

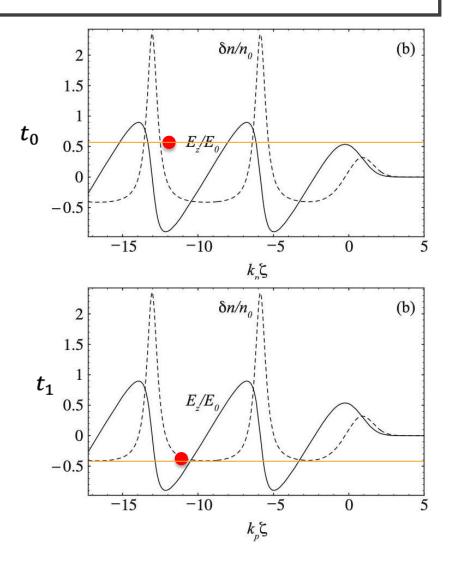


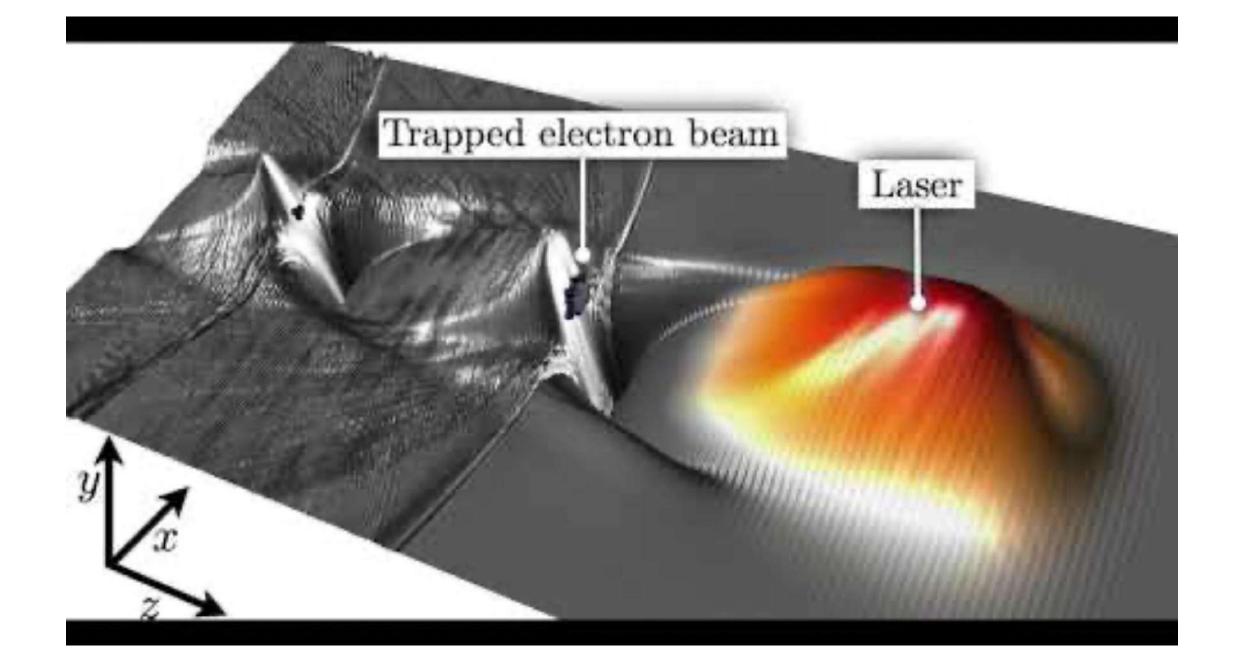
What Must LWFA Overcome?

Dephasing and Pump Depletion

- As the electrons accelerate, their group velocity eventually outruns the group velocity of the laser
- Must turn off plasma source once this switchover occurs
- Laser runs out of energy overtime

$$v_g = \sqrt{1 - \frac{n_e}{n_c}}$$





THANKS FOR LISTENING!